

**Tensors, field strength and Lorentz transform of electromagnetic fields****1 Symmetric and antisymmetric tensors****1.1 Symmetrization, antisymmetrization and covariance**

Consider a quadridimensional 2-tensor  $M^{\mu\nu}$ .

1. Provide a separation of  $M^{\mu\nu}$  as a sum of two symmetric ( $S$ ) and antisymmetric ( $A$ ) tensors:

$$M^{\mu\nu} = S^{\mu\nu} + A^{\mu\nu}. \quad (1)$$

2. Recall the way  $M^{\mu\nu}$  transform under a arbitrary Lorentz transformation  $\Lambda$ , encoded through its matrix elements  $\Lambda^\rho_\sigma$ .

3. Show that the decomposition (1) is covariant under Lorentz transformations.

4. Write the Lorentz transformation of  $M^{\mu\nu}$  in a way which makes this separation explicit.

**1.2 Transformation under boosts**

We now consider a Lorentz boost of a frame  $F$  to a frame  $F'$  along the  $x$  axis, encoded by  $\beta = v/c$  and  $\gamma = 1/\sqrt{1 - \beta^2}$ .

5. Recall the explicit expression of  $\Lambda$ .
6. We focus on the case of a symmetric tensor  $S^{\mu\nu}$ .

Write the Lorentz transformation of the various components of  $S$  under the above boost.

7. We focus on the case of an antisymmetric tensor  $A^{\mu\nu}$ .

- a. What can be said about  $A^{00}, A^{11}, A^{22}, A^{33}$  and their transformations?

- b. How does  $A^{23}$  transforms?

- c. Compare the transformation of  $A^{12}, A^{13}$  and  $A^{02}, A^{03}$  with the transformation of  $x^1$  and  $x^0$ . Deduce the transformation of these components.

d. Show that  $A^{01}$  is invariant under these boosts.

## 2 Lorentz transformations of electromagnetic fields

We now apply the previous results to the field strength

$$F^{\mu\nu} = \begin{pmatrix} 0 & -E^1 & -E^2 & -E^3 \\ E^1 & 0 & -B^3 & B^2 \\ E^2 & B^3 & 0 & -B^1 \\ E^3 & -B^2 & B^1 & 0 \end{pmatrix},$$

where  $\vec{E}$  and  $\vec{B}$  are the electric and magnetic fields. We use a system of units such that  $c = 1$ .

8. Show that under the considered boost along  $x$ , these fields transform as

$$E'^1 = E^1, \quad (2)$$

$$E'^2 = \gamma(E^2 - \beta B^3), \quad (3)$$

$$E'^3 = \gamma(E^3 + \beta B^2) \quad (4)$$

and

$$B'^1 = B^1, \quad (5)$$

$$B'^2 = \gamma(B^2 + \beta E^3), \quad (6)$$

$$B'^3 = \gamma(B^3 - \beta E^2). \quad (7)$$

9. Consider an arbitrary boost along the direction  $\vec{n}$  ( $\vec{n}^2 = 1$ ), i.e. with a velocity  $\vec{v} = \beta\vec{n}$ .

Show that under such a boost,

$$\vec{E}' = (\vec{E} \cdot \vec{n})\vec{n} + \gamma \left[ \vec{E} - (\vec{E} \cdot \vec{n})\vec{n} \right] + \gamma \vec{v} \wedge \vec{B}, \quad (8)$$

$$\vec{B}' = (\vec{B} \cdot \vec{n})\vec{n} + \gamma \left[ \vec{B} - (\vec{B} \cdot \vec{n})\vec{n} \right] - \gamma \vec{v} \wedge \vec{E}. \quad (9)$$

10. Show that this can be rewritten in the form

$$\vec{E}' = (\vec{E} \cdot \vec{n})\vec{n} + \gamma \left[ \vec{n} \wedge (\vec{E} \wedge \vec{n}) + \vec{v} \wedge \vec{B} \right], \quad (10)$$

$$\vec{B}' = (\vec{B} \cdot \vec{n})\vec{n} + \gamma \left[ \vec{n} \wedge (\vec{B} \wedge \vec{n}) - \vec{v} \wedge \vec{E} \right]. \quad (11)$$

*Hint:* use  $\vec{a} \wedge (\vec{b} \wedge \vec{c}) = (\vec{a} \cdot \vec{c})\vec{b} - (\vec{a} \cdot \vec{b})\vec{c}$ .

11. In the non-relativistic limit, simplify the transformations (8) and (9), keeping linear terms in  $\beta$ .