## M2/CFP/Parcours of Physique Théorique

Invariances in physics and group theory

## Exercises on Cartan algebra, roots, weights...

Equations refer to the lecture of Jean-Bernard Zuber, Chapter 3, version 2013.
A. Cartan algebra and roots

1. Show that any element $X$ of $\mathfrak{g}$ may be written as $X=\sum x^{i} H_{i}+\sum_{\alpha \in \Delta} x^{\alpha} E_{\alpha}$ with the notations of § 3.1.2.

For an arbitrary $H$ in the Cartan algebra, determine the action of $\operatorname{ad}(H)$ on such a vector $X$; conclude that $\operatorname{ad}(H) \operatorname{ad}\left(H^{\prime}\right) X=\sum_{\alpha \in \Delta} x^{\alpha} \alpha(H) \alpha\left(H^{\prime}\right) E_{\alpha}$ and taking into account that the eigenspace of each root $\alpha$ has dimension 1 , cf point $(*)$ 2. of § 3.1.2, that the Killing form reads

$$
\begin{equation*}
\left(H, H^{\prime}\right)=\operatorname{tr}\left(\operatorname{ad}(\mathrm{H}) \operatorname{ad}\left(\mathrm{H}^{\prime}\right)\right)=\sum_{\alpha \in \Delta} \alpha(\mathrm{H}) \alpha\left(\mathrm{H}^{\prime}\right) . \tag{1}
\end{equation*}
$$

[Solution : For a fixed value of $j, H_{j}$ has $\operatorname{dim} \mathfrak{g}$ independent eigenvectors (since $H_{j}$ acts on $\mathfrak{g}$ ). These eigenvectors thus provide a set of generators, allowing to write

$$
\begin{equation*}
\forall X \in \mathfrak{g}, X=\sum x^{i} H_{i}+\sum_{\alpha \in \Delta} x^{\alpha} E_{\alpha} \tag{2}
\end{equation*}
$$

For $H \in \mathfrak{h}$,

$$
\operatorname{ad} H X=[H, X]=\sum_{\alpha \in \Delta} x^{\alpha}\left[H, E_{\alpha}\right]=\sum_{\alpha \in \Delta} x^{\alpha} \alpha(H) E_{\alpha} .
$$

Thus

$$
\operatorname{ad}(H) \operatorname{ad}\left(H^{\prime}\right) X=\sum_{\alpha \in \Delta} x^{\alpha} \alpha\left(H^{\prime}\right)\left[H, E_{\alpha}\right]=\sum_{\alpha \in \Delta} x^{\alpha} \alpha\left(H^{\prime}\right) \alpha(H) E_{\alpha}
$$

and

$$
\alpha\left(H^{\prime}\right) \alpha(H) E_{\alpha}=\alpha\left(H^{\prime}\right) \alpha(H) E_{\alpha}
$$

for a given $\alpha$. Each root $\alpha$ corresponds to an eigenspace spanned by $E_{\alpha}$, of dimension 1. Thus

$$
\operatorname{tr}\left[\operatorname{ad}(\mathrm{H}) \operatorname{ad}\left(\mathrm{H}^{\prime}\right)\right]=\sum_{\alpha \in \Delta} \alpha\left(\mathrm{H}^{\prime}\right) \alpha(\mathrm{H}) .
$$

]
2. One wants to show that roots $\alpha$ defined by (3.5) or (3.6) generate all the dual space $\mathfrak{h}^{*}$ of the Cartan subalgebra $\mathfrak{h}$. Prove that if it were not so, there would exist an element $H$ of $\mathfrak{h}$ such that

$$
\begin{equation*}
\forall \alpha \in \Delta \quad \alpha(H)=0 . \tag{3}
\end{equation*}
$$

Using (1) show that this would imply $\forall H^{\prime} \in \mathfrak{h},\left(H, H^{\prime}\right)=0$. Why is that impossible in a semi-simple algebra? (see the discussion before equation (3.10)).
[Solution: One is looking for $H=\sum_{i} h^{i} H_{i} \in \mathfrak{h}$, so that $\alpha(H)=\sum_{i} h^{i} \alpha_{(i)}=0$ for every $\alpha \in \Delta$. Thus, $H$ should obey $|\Delta|$ equations, which would lead to a number of undetermined parameters larger or equal to $\operatorname{dim} \mathfrak{h}-|\Delta|,\left\{h^{i}, i \in\{1 \cdots \operatorname{dim} \mathfrak{h}\}\right\}$ being a solution of this set of equations. In the case where $|\Delta|<\operatorname{dim} \mathfrak{h}^{*}=\operatorname{dim} \mathfrak{h}$, a non-trivial solution (i.e. such that $h^{i}$ do not all vanish) thus exists, and $H$ is non zero.

The expression

$$
\left(H, H^{\prime}\right)=\sum_{\alpha \in \Delta} \alpha\left(H^{\prime}\right) \alpha(H)
$$

leads to

$$
\forall H^{\prime} \in \mathfrak{h},\left(H, H^{\prime}\right)=0 .
$$

Since $\left(H, E_{\alpha}\right)=0$, we thus deduce that

$$
\forall X \in \mathfrak{h},(H, X)=0
$$

thus showing that the Killing form is degenerate on $\mathfrak{h}$, in contradiction with the semi-simplicity of $\mathfrak{h}$. Thus, $\Delta$ spans $\mathfrak{h}^{*}$ (and thus $|\Delta| \geq \operatorname{dim} \mathfrak{h}^{*}=\operatorname{dim} h$ ).]
3. Variant of the previous argument : under the assumption of 2 . and thus of (3), show that $H$ would commute with all $H_{i}$ and all the $E_{\alpha}$, thus would belong to the center of $\mathfrak{g}$. Prove that the center of an algebra is an abelian ideal. Conclude in the case of a semi-simple algebra.
[Solution : Assume that $\exists H \in \mathfrak{h}$ such that $\forall \alpha \in \Delta, \alpha(H)=0$. Then, on one hand $\left[H, H_{i}\right]=0$, while on the other hand $\forall \alpha \in \Delta,\left[H, E_{\alpha}\right]=\alpha(H) E_{\alpha}=0$ so that $H \in Z_{\mathfrak{g}}$. Now, $X \in Z_{\mathfrak{g}} \Leftrightarrow \forall Y \in \mathfrak{g},[X, Y]=0$. Obviously, $\left[Z_{\mathfrak{g}}, \mathfrak{g}\right] \subset Z_{\mathfrak{g}}$ since $\left[Z_{\mathfrak{g}}, \mathfrak{g}\right]=\{0\}$, showing that $Z_{\mathfrak{g}}$ is an abelian ideal of $\mathfrak{g}$ (abelian since every element of $Z_{\mathfrak{g}}$ commute with every element of $\mathfrak{g}$, in particular those of $Z_{\mathfrak{g}}$ ). Based on the semi-simplicity of $\mathfrak{g}$, this implies that $Z_{\mathfrak{g}}=\{0\}$, proving that $H \in Z_{\mathfrak{g}}$ is impossible.]
B. Computation of the $N_{\alpha \beta}$

1. Show that the real constants $N_{\alpha \beta}$ satisfy $N_{\alpha \beta}=-N_{\beta \alpha}$ and, by complex conjugation of $\left[E_{\alpha}, E_{\beta}\right]=N_{\alpha \beta} E_{\alpha+\beta}$ that

$$
\begin{equation*}
N_{\alpha, \beta}=-N_{-\alpha,-\beta} . \tag{4}
\end{equation*}
$$

[Solution : $\left[E_{\alpha}, E_{\beta}\right]=N_{\alpha, \beta} E_{\alpha+\beta}$ if $\alpha+\beta$ is a root. Since $\left[E_{\alpha}, E_{\beta}\right]=-\left[E_{\beta}, E_{\alpha}\right]$, we thus have $N_{\beta, \alpha}=-N_{\alpha, \beta}$.
$\left[H, E_{\alpha}\right]=\alpha E_{\alpha}$ leads to $\left[H, E_{\alpha}\right]^{\dagger}=-\left[H, E_{\alpha}^{\dagger}\right]=\alpha E_{\alpha}^{\dagger}$ since $H$ is hermitian. Thus, $\left[H, E_{\alpha}^{\dagger}\right]=-\alpha E_{\alpha}^{\dagger}$ and thus $E_{\alpha}^{\dagger}=E_{-\alpha}$. From $\left[E_{\alpha}, E_{\beta}\right]=N_{\alpha, \beta} E_{\alpha+\beta}$ we thus have $\left[E_{\alpha}, E_{\beta}\right]^{\dagger}=N_{\alpha, \beta} E_{\alpha+\beta}^{\dagger}=N_{\alpha, \beta} E_{-\alpha-\beta}$. This leads to

$$
N_{\alpha, \beta}=N_{-\beta,-\alpha}=-N_{-\alpha,-\beta}=-N_{\beta, \alpha} .
$$

]
2. Consider three roots satisfying $\alpha+\beta+\gamma=0$. Writing the Jacobi identity for the triplet $E_{\alpha}, E_{\beta}, E_{\gamma}$, show that $\alpha_{(i)} N_{\beta \gamma}+$ cycl. $=0$. Derive from it the relation

$$
\begin{equation*}
N_{\alpha \beta}=N_{\beta,-\alpha-\beta}=N_{-\alpha-\beta, \alpha} . \tag{5}
\end{equation*}
$$

[Solution: The Jacobi identity reads

$$
\left[\left[E_{\alpha}, E_{\beta}\right], E_{\gamma}\right]+\left[\left[E_{\beta}, E_{\gamma}\right], E_{\alpha}\right]+\left[\left[E_{\gamma}, E_{\alpha}\right], E_{\beta}\right]=0
$$

and thus

$$
N_{\alpha, \beta}\left[E_{\alpha+\beta}, E_{\gamma}\right]+N_{\beta, \gamma}\left[E_{\beta+\gamma}, E_{\alpha}\right]+N_{\gamma, \alpha}\left[E_{\alpha+\gamma}, E_{\beta}\right]=0,
$$

which can be written as

$$
N_{\alpha, \beta}\left[E_{-\gamma}, E_{\gamma}\right]+N_{\beta, \gamma}\left[E_{-\alpha}, E_{\alpha}\right]+N_{\gamma, \alpha}\left[E_{-\beta}, E_{\beta}\right]=0
$$

i.e.

$$
N_{\alpha, \beta} H_{-\gamma}+N_{\beta, \gamma} H_{-\alpha}+N_{\gamma, \alpha} H_{-\beta}=0,
$$

so that

$$
N_{\beta, \gamma}\left(-\alpha_{(j)} H_{j}\right)+\text { cycl.perm. }=0 .
$$

Calculating ( $H_{i}$, ) for a fixed value of $i$, we thus get

$$
\alpha_{(i)} N_{\beta, \gamma}+\text { cycl.perm. }=0 .
$$

We thus have

$$
\alpha N_{\beta, \gamma}+\beta N_{\gamma, \alpha}+\gamma N_{\alpha, \beta}=0
$$

so that

$$
\alpha\left(N_{\beta,-\alpha-\beta}-N_{\alpha, \beta}\right)+\beta\left(N_{-\alpha-\beta, \alpha}-N_{\alpha, \beta}\right)=0 .
$$

This identity is valid for any set of roots $\alpha, \beta, \gamma$ satisfying $\alpha+\beta+\gamma=0$, so that $\alpha$ and $\beta$ are linearly independent, except if $\beta= \pm \alpha$. Now :

- If $\alpha=\beta$, then $\gamma=-2 \alpha$ which is impossible since the only roots of the form $\lambda \alpha$ are $\pm \alpha$.
- If $\alpha=-\beta$, then $\gamma=0$ which is again impossible.

We thus conclude, based on the linear independence of $\alpha$ and $\beta$, that

$$
N_{\beta,-\alpha-\beta}=N_{\alpha, \beta}=N_{-\alpha-\beta, \alpha} .
$$

This identity remains trivially satisfied for $\beta=\alpha$, since on one hand $\left[E_{\alpha}, E_{\alpha}\right]=0$, i.e. $N_{\alpha, \alpha}=0$ and on the other hand $N_{\alpha,-2 \alpha}=0$ since $-2 \alpha$ is not a root. Similarly, it is also valid in the case $\beta=-\alpha$ since $N_{\alpha,-\alpha}=N_{-\alpha, 0}=N_{0, \alpha}=0$. The above identity is thus valid for any set of roots $\alpha$ and $\beta$.]
3. Considering the $\alpha$-chain through $\beta$ and the two integers $p$ et $q$ defined in $\S$ 3.2.1, write the Jacobi identity for $E_{\alpha}, E_{-\alpha}$ and $E_{\beta+k \alpha}$, with $p \leq k \leq q$, and show that it implies

$$
\langle\alpha, \beta+k \alpha\rangle=N_{\alpha, \beta+k \alpha} N_{\alpha, \beta+(k-1) \alpha}+N_{\beta+k \alpha, \alpha} N_{-\alpha, \beta+(k+1) \alpha} .
$$

Let $f(k):=N_{\alpha, \beta+k \alpha} N_{-\alpha,-\beta-k \alpha}$. Using the relations (5), show that the previous equation may be recast as

$$
\begin{equation*}
\langle\alpha, \beta+k \alpha\rangle=f(k)-f(k-1) . \tag{6}
\end{equation*}
$$

[Solution : The Jacobi identity reads

$$
\left.\left[\left[E_{\alpha}, E_{-\alpha}\right], E_{\beta+k \alpha}\right]+\left[\left[E_{-\alpha}, E_{\beta+k \alpha}\right], E_{\alpha}\right]+\left[\left[E_{\beta+k \alpha}\right], E_{\alpha}\right], E_{-\alpha}\right]=0
$$

i.e.

$$
\left[H_{\alpha}, E_{\beta+k \alpha}\right]+N_{-\alpha, \beta+k \alpha}\left[E_{\beta+(k-1) \alpha}, E_{\alpha}\right]+N_{\beta+k \alpha, \alpha}\left[E_{\beta+(k+1) \alpha}, E_{-\alpha}\right]=0,
$$

and thus

$$
\begin{aligned}
& \sum_{i} \alpha(i)[\beta(i)+k \alpha(i)] E_{\beta+k \alpha}+N_{-\alpha, \beta+k \alpha} N_{\beta+(k-1) \alpha, \alpha} E_{\beta+k \alpha} \\
& \quad+N_{\beta+k \alpha, \alpha} N_{\beta+(k+1) \alpha,-\alpha} E_{\beta+k \alpha}=0
\end{aligned}
$$

From $\sum_{i} \alpha(i)[\beta(i)+k \alpha(i)]=\langle\alpha, \beta+k \alpha\rangle, N_{\beta+(k-1) \alpha, \alpha}=-N_{\alpha, \beta+(k-1) \alpha}$ and $N_{\beta+(k+1) \alpha,-\alpha}=$ $-N_{-\alpha, \beta+(k+1) \alpha}$ we get, since $E_{\beta+k \alpha} \neq 0$,

$$
\langle\alpha, \beta+k \alpha\rangle=N_{-\alpha, \beta+k \alpha} N_{\alpha, \beta+(k-1) \alpha}+N_{\beta+k \alpha, \alpha} N_{-\alpha, \beta+(k+1) \alpha} .
$$

From the definition of $f$, we have
$f(k)=N_{\alpha, \beta+k \alpha} N_{-\alpha,-\beta-k \alpha}=-N_{\alpha, \beta+k \alpha}^{2}=N_{\beta+k \alpha, \alpha} N_{-\beta-k \alpha,-\alpha}=N_{\beta+k \alpha, \alpha} N_{-\alpha, \beta+(k+1) \alpha}$.
Thus,

$$
f(k-1)=N_{\alpha, \beta+(k-1) \alpha} N_{-\alpha,-\beta-(k-1) \alpha} .
$$

Now, since $N_{-\alpha,-\beta-(k-1) \alpha}=N_{\beta+k \alpha,-\alpha}=-N_{-\alpha, \beta+k \alpha}$, where we have used $N_{\gamma, \delta}=$ $-N_{-\gamma-\delta, \gamma}$, we thus have

$$
\langle\alpha, \beta+k \alpha\rangle=f(k)-f(k-1) .
$$

]
4. What are $f(q)$ et $f(q-1)$ ? Show that the recursion relation (6) is solved by

$$
\begin{equation*}
f(k)=-\left(N_{\alpha, \beta+k \alpha}\right)^{2}=(k-q)\left\langle\alpha, \beta+\frac{1}{2}(k+q+1)\right\rangle . \tag{7}
\end{equation*}
$$

What is $f(p-1)$ ? Show that the expression (3.21) is recovered. Show that (7) is in accord with (3.23). The sign of the square root is still to be determined ..., see [Gi].
[Solution : The definition of $f$ leads to $f(q)=N_{\alpha, \beta+q \alpha} N_{-\alpha,-\beta-q \alpha}$. Since $\left[E_{\alpha}, E_{\beta+q \alpha}\right]=$ $N_{\alpha, \beta+q \alpha} E_{\beta+(q+1) \alpha}$ while $\beta+(q+1) \alpha$ is not root, as implied by the definition of $q$, we obtain $N_{\alpha, \beta+q \alpha}=0$ and thus

$$
f(q)=0 \text {. }
$$

From the definition of $f$,

$$
f(q-1)=-N_{\alpha, \beta+(q-1) \alpha}^{2} .
$$

On the other hand, $\langle\alpha, \beta+q \alpha\rangle=f(q)-f(q-1)$, so that

$$
f(q-1)=-\langle\alpha, \beta+q \alpha\rangle=-N_{\alpha, \beta+(q-1) \alpha}^{2} .
$$

From the expression $\langle\alpha, \beta\rangle+k\langle\alpha, \alpha\rangle=f(k)-f(k-1)$ we get
$f(q)-f(q-1)+f(q-1)-f(q-2)+\cdots+f(k+1)-f(k)=-f(k)=(q-k)\langle\alpha, \beta\rangle+\sum_{m=k+1}^{q} m\langle\alpha, \alpha\rangle$.
From $\sum_{m=k+1}^{q} m=-\sum_{m=1}^{k} m+\sum_{m=1}^{q}=\frac{q(q+1)}{2}-\frac{k(k+1)}{2}=-\frac{k-q}{2}(1+k+q)$ we get

$$
f(k)=(k-q)\langle\alpha, \beta\rangle+\frac{1}{2}\langle\alpha, \alpha\rangle(1+k+q)(k-q)
$$

so that

$$
f(k)=(k-q)\left\langle\alpha, \beta+\frac{1}{2}(k+q+1) \alpha\right\rangle=-N_{\alpha, \beta+k \alpha}^{2} .
$$

We now compute

$$
f(p-1)=N_{\alpha, \beta+(p-1) \alpha} N_{-\alpha,-\beta-(p-1) \alpha} .
$$

The second coefficient reads $N_{-\alpha,-\beta-(p-1) \alpha}=N_{\beta+p \alpha,-\alpha}=-N_{-\alpha, \beta+p \alpha}$. Since

$$
\left[E_{-\alpha}, E_{\beta+p \alpha}\right]=N_{-\alpha, \beta+p \alpha} E_{\beta+(p-1) \alpha}
$$

and using the fact that $\beta+(p-1) \alpha$ is not a root by definition of $q$, we deduce that $N_{-\alpha, \beta+p \alpha}=0$ and thus $f(p-1)=0$. Therefore,

$$
f(p)=\langle\alpha, \beta+p \alpha\rangle=-N_{\alpha, \beta+p \alpha}^{2}=(p-q)\left\langle\alpha, \beta+\frac{1}{2}(p+q+1) \alpha\right\rangle,
$$

so that

$$
\langle\alpha, \beta\rangle(p-q-1)=\left[p-\frac{1}{2}(p-q)(p+q+1)\right]\langle\alpha, \alpha\rangle=-\frac{1}{2}(p-q-1)(p+q)\langle\alpha, \alpha\rangle .
$$

Therefore,

$$
2\langle\alpha, \beta\rangle=(-p-q)\langle\alpha, \alpha\rangle=m\langle\alpha, \alpha\rangle .
$$

We thus have

$$
2 \frac{\langle\alpha, \beta\rangle}{\langle\alpha, \alpha\rangle}=m \in \mathbb{Z}
$$

We finally get
$f(0)=-N_{\alpha, \beta}^{2}=-q\left\langle\alpha, \beta+\frac{1}{2}(q+1) \alpha\right\rangle=-q\left[\langle\alpha, \beta\rangle+\frac{1}{2}(q+1)\langle\alpha, \alpha\rangle\right]=-\frac{q(1-p)}{2}\langle\alpha, \alpha\rangle$ and thus

$$
\left|N_{\alpha, \beta}\right|=\sqrt{\frac{q(1-p)}{2}\langle\alpha, \alpha\rangle} .
$$

]
C. Study of the $B_{l}=s o(2 l+1)$ and $G_{2}$ algebras

1. $s o(2 l+1)=B_{l}, l \geq 2$
a. What is the dimension of the group $\mathrm{SO}(2 l+1)$ or of its Lie algebra $\operatorname{so}(2 l+1)$ ?
[Solution : We start from the defining relation

$$
O O^{t}=O^{t} O=1
$$

with $O$ expanded as $O=1+i T$, i.e. $(1+i T)\left(1+i T^{t}\right)=1$ and thus $i\left(T+T^{t}\right)=0$, showing that $T$ is antisymmetric.

If $T$ is an $n \times n$ matrix (for the case of $O(n)$ or $S O(n)$ ), $T$ is characterized by $\left(n^{2}-n\right) / 2$ elements, so that for $n=2 \ell+1$,

$$
\operatorname{dim} o(2 \ell+1)=\ell(2 \ell+1) .
$$

This dimension is the same for so $(2 \ell+1)$ since $\operatorname{det} O=1$ does not introduce any additional constraint on the Lie algebra (it only selects the connected component of the identity when considering the group).]
b. What is the rank of the algebra? (Hint : diagonalize a matrix of so $(2 l+1)$ on $\mathbb{C}$, or write it as a diagonal of real $2 \times 2$ blocks, see lecture notes, $\S 3.1$ )
[Solution : Since $O$ is diagonalizable in the form

$$
D=\operatorname{diag}\left(0,\left(\begin{array}{cc}
0 & \mu_{j} \\
-\mu_{j} & 0
\end{array}\right)_{j=1, \cdots, \ell}\right)
$$

with real $\mu_{j}$ (see chapter 3), we thus get

$$
\operatorname{rank} s o(2 \ell+1)=\ell .
$$

]
c. How many roots does the algebra have? How many positive roots? How many simple?
[Solution : \# of roots = dimension $-\operatorname{rank}=\ell(2 \ell+1)-\ell=2 l^{2}$.
Among these roots, half of them are positive, i.e. $\ell^{2}$. The number of simple roots is equal to the rank, i.e. $\ell$.]
d. Let $e_{i}, i=1, \cdots, l$ be a orthonormal basis in $\mathbb{R}^{l},\left\langle e_{i}, e_{j}\right\rangle=\delta_{i j}$. Consider the set of vectors

$$
\Delta=\left\{ \pm e_{i}, 1 \leq i \leq l\right\} \cup\left\{ \pm e_{i} \pm e_{j}, \quad 1 \leq i<j \leq l\right\}
$$

What is the cardinal of $\Delta$ ?
[Solution : Card $\left\{ \pm e_{i}, 1 \leq i \leq l\right\}=2 \ell$ and $\operatorname{Card}\left\{ \pm e_{i} \pm e_{j}, \quad 1 \leq i<j \leq l\right\}=$ $4 C_{\ell}^{2}=2 l(l-1)$, therefore Card $\Delta=2 l^{2}$.]
$\Delta$ is the set of roots of the algebra so $(2 l+1)$.
e. A basis of simple roots is given $\alpha_{i}=e_{i}-e_{i+1}, i=1, \cdots, l-1$, et $\alpha_{l}=e_{l}$. Explain why the roots

$$
\begin{align*}
e_{i} & =\sum_{i \leq k \leq l} \alpha_{k}, \quad 1 \leq i \leq l, \\
e_{i}-e_{j} & =\sum_{i \leq k<j} \alpha_{k}, \quad 1 \leq i<j \leq l,  \tag{8}\\
e_{i}+e_{j} & =\sum_{i \leq k<j} \alpha_{k}+2 \sum_{j \leq k \leq l} \alpha_{k}, \quad 1 \leq i<j \leq l,
\end{align*}
$$

qualify as positive roots.
[Solution : There are

$$
\begin{align*}
& e_{i}=\sum_{i \leq k \leq l} \alpha_{k}, \quad 1 \leq i \leq l, \quad \text { Card }=\ell, \\
& e_{i}-e_{j}=\sum_{i \leq k<j} \alpha_{k}, \quad 1 \leq i<j \leq l, \quad \operatorname{Card}=\frac{\ell(\ell-1)}{2},  \tag{9}\\
& e_{i}+e_{j}=\sum_{i \leq k<j} \alpha_{k}+2 \sum_{j \leq k \leq l} \alpha_{k}, \quad 1 \leq i<j \leq l, \quad \operatorname{Card}=\frac{\ell(\ell-1)}{2},
\end{align*}
$$

there are thus $\left|\Delta_{+}\right|=l^{2}$ of such roots, which expand on the set of simple roots with integer coefficients. Including their opposite (negative roots), they correctly reproduce the whole set $\Delta$.] Check that assertion on the case of $B_{2}=s o(5)$.
[Solution: For $B_{2}=s o(5), \Delta=\left\{ \pm e_{1}, \pm e_{2}, \pm e_{1} \pm e_{2}\right\}$ with Card $\Delta=8$. The $\ell=2$ (which is equal to the rank of $B_{2}=s o(5)$ ) simple roots are $\alpha_{1}=e_{1}-e_{2}$ and $\alpha_{2}=e_{2}$. And $\Delta_{+}=\left\{e_{1}, e_{2}, e_{1}-e_{2}, e_{1}+e_{2}\right\}$ with Card $\left.\Delta_{+}=4.\right]$
f. Compute the Cartan matrix and check that it agrees with the Dynkin diagram given in the notes.
[Solution: Obviously, $2 \frac{\left\langle\alpha_{i}, \alpha_{i}\right\rangle}{\left\langle\alpha_{i}, \alpha_{i}\right\rangle}=2$. Then,

$$
2 \frac{\left\langle\alpha_{j+1}, \alpha_{j}\right\rangle}{\left\langle\alpha_{j}, \alpha_{j}\right\rangle}=2 \frac{\left\langle e_{j+1}-e_{j+2}, e_{j}-e_{j+1}\right\rangle}{\left\langle e_{j}-e_{j+1}, e_{j}-e_{j+1}\right\rangle}=\frac{2(-1)}{2}=-1
$$

and

$$
2 \frac{\left\langle\alpha_{j}, \alpha_{j+1}\right\rangle}{\left\langle\alpha_{j+1}, \alpha_{j+1}\right\rangle}=-1 .
$$

Finally,

$$
2 \frac{\left\langle\alpha_{\ell-1}, \alpha_{\ell}\right\rangle}{\left\langle\alpha_{\ell}, \alpha_{\ell}\right\rangle}=2 \frac{\left\langle e_{\ell-1}-e_{\ell}, e_{\ell}\right\rangle}{\left\langle e_{\ell}, e_{\ell}\right\rangle}=-2,
$$

and

$$
2 \frac{\left\langle\alpha_{\ell}, \alpha_{\ell-1}\right\rangle}{\left\langle\alpha_{\ell-1}, \alpha_{\ell-1}\right\rangle}=2 \frac{\left\langle e_{\ell}, e_{\ell-1}-e_{\ell}\right\rangle}{\left\langle e_{\ell-1}-e_{\ell}, e_{\ell-1}-e_{\ell}\right\rangle}=\frac{2(-1)}{2}=-1,
$$

which can be summarized by

$$
C_{i j}=2 \frac{\left\langle\alpha_{i}, \alpha_{j}\right\rangle}{\left\langle\alpha_{j}, \alpha_{j}\right\rangle}=\left\{\begin{aligned}
2 & \text { si } 1 \leq i=j \leq l \\
-1 & \text { si } 1 \leq i=(j \pm 1) \leq l-1 \\
-2 & \text { si } i=l-1, j=l \\
-1 & \text { si } i=l, j=l-1
\end{aligned}\right.
$$

corresponding to the Dynkin diagram illustrated in Fig. 1. The roots of $B_{2}$ are


Figure 1 - The Dynkin diagram for the algebra $B_{\ell}$.
illustrated in Fig. 2 ]
g. Compute the sum $\rho$ of positive roots.
[Solution :

$$
\begin{gathered}
2 \rho=\sum_{i=1}^{\ell} e_{i}+\sum_{1 \leq i<j \leq \ell}\left(e_{i}-e_{j}+e_{i}+e_{j}\right)=\sum_{i=1}^{\ell} e_{i}+2 \sum_{1 \leq i<j \leq \ell} e_{i}=\sum_{1 \leq i<j \leq \ell}[1+2(\ell-i)] e_{i} \\
=(2 l-1) e_{1}+(2 l-3) e_{2}+\cdots+(2 l-2 i+1) e_{i}+\cdots+3 e_{l-1}+e_{l} \\
=(2 l-1)\left(\alpha_{1}+\cdots \alpha_{\ell}\right)+(2 l-3)\left(\alpha_{2}+\cdots \alpha_{\ell}\right)+\cdots+(2 l-2 i+1)\left(\alpha_{i}+\cdots \alpha_{\ell}\right) \\
+\cdots+3\left(\alpha_{\ell-1}+\alpha_{\ell}\right)+\alpha_{l} \\
=(2 l-1) \alpha_{1}+2(2 l-2) \alpha_{2}+\cdots+i(2 l-i) \alpha_{i}+\cdots+l^{2} \alpha_{l} .
\end{gathered}
$$



Figure 2 - The roots of $B_{2}$. In red, the two simple roots $\alpha_{1}$ and $\alpha_{2}$. In black, the two other positive roots, and in cyan the four negative roots.
]
h. The Weyl is the ("semi-direct") product $W \equiv \mathcal{S}_{l} \ltimes\left(\mathbb{Z}_{2}\right)^{l}$, which acts on the $e_{i}$ (hence on weights and roots) by permutation and by independent sign changes $e_{i} \mapsto( \pm 1)_{i} e_{i}$. What is its order? In the case of $B_{2}$, check that assertion and draw the first Weyl chamber.
[Solution: $\left|W_{\ell}\right|=2^{l} . l$ !, and for $B_{2},\left|W_{2}\right|=8$. Permutations and change of sign of $e_{1}$ and $e_{2}$ indeed corresponds to (products of) symmetries in "planes" orthogonal to $\alpha_{1}$ and $\alpha_{2}$; they do not modify the picture of Fig. 2. It is worth obtaining $W_{2}$ in a pedestrian way. We easily get

$$
\left\{\begin{array}{l}
\alpha_{1}^{\perp} e_{1}=e_{2}, \quad \alpha_{2}^{\perp} e_{1}=e_{1} \\
\alpha_{1}^{\perp} e_{2}=e_{1}, \quad \alpha_{2}^{\perp} e_{2}=-e_{2} .
\end{array}\right.
$$

Thus

$$
\left\{\begin{array}{l}
\alpha_{1}^{\perp} \alpha_{2}^{\perp} e_{1}=e_{2}, \quad \alpha_{2}^{\perp} \alpha_{1}^{\perp} e_{1}=-e_{2}, \\
\alpha_{1}^{\perp} \alpha_{2}^{\perp} e_{2}=-e_{1}, \quad \alpha_{2}^{\perp} \alpha_{1}^{\perp} e_{2}=e_{1} .
\end{array}\right.
$$

This shows that $\alpha_{1}^{\perp} \alpha_{2}^{\perp}=-\alpha_{2}^{\perp} \alpha_{1}^{\perp}$ is a rotation of angle $-\pi / 2$ (in the oriented plane $e_{1}, e_{2}$ ). Thus, any word made of $\mathbb{I},-\mathbb{I}, \alpha_{1}^{\perp}, \alpha_{2}^{\perp}$ reduces to the elements of $W_{2}=\left\{\mathbb{I}, \alpha_{1}^{\perp}, \alpha_{2}^{\perp}, \alpha_{1}^{\perp} \alpha_{2}^{\perp},-\mathbb{I},-\alpha_{1}^{\perp},-\alpha_{2}^{\perp},-\alpha_{1}^{\perp} \alpha_{2}^{\perp}\right\}$. Equivalently, the Weyl group is generated by the set of reflections in the hyperplane (here line) orthogonal to the various $\alpha$ 's. The first Weyl chamber $\mathcal{C}_{1}=\left\{\lambda \mid\left\langle\lambda, \alpha_{i}\right\rangle \geq 0\right\}$ is the octant between $\alpha_{1}+\alpha_{2}$ and $\alpha_{1}+2 \alpha_{2}$.]
i. Show that the vectors $\Lambda_{i}=\sum_{j=1}^{i} e_{j}, i=1, \cdots, l-1, \Lambda_{l}=\frac{1}{2} \sum_{j=1}^{l} e_{j}$ are the fundamental weights.
[Solution: The fundamental weights are defined in such a way that $\Lambda_{i}$ is orthogonal


Figure 3 - The weight diagram for the algebra $B_{2}$. The two simple roots $\alpha_{1}$ and $\alpha_{2}$ are displayed as red lines. The two fundamental weights $\Lambda_{1}$ and $\Lambda_{2}$ are displayed as blue lines. The lattice of roots is shown as red dots, while the lattice of weights is shown as blue dots. The first Weyl chamber $\mathcal{C}_{1}$ is shown as a dark-grey zone. The other Weyl chambers, shown as grey regions of various intensities, are obtained by acting on $\mathcal{C}_{1}$ with the various element of the Weyl group, as explicitely indicated.
to every root simple root $\alpha_{j}(j \neq i)$, its normalization being fixed by $2 \frac{\left\langle\alpha_{i}, \Lambda_{i}\right\rangle}{\left\langle\alpha_{i}, \alpha_{i}\right\rangle}=1$. The denominators read

$$
\begin{aligned}
& \left\langle\alpha_{i}, \alpha_{i}\right\rangle=\left\langle e_{i}-e_{i+1}, e_{i}-e_{i+1}\right\rangle=2 \quad \text { for } i=1, \cdots, \ell-1 \\
& \left\langle\alpha_{\ell}, \alpha_{\ell}\right\rangle=1
\end{aligned}
$$

The construction can be easily performed by recursion. First, $\Lambda_{1}$ should be orthogonal to $\alpha_{i}(i>1)$, so that, from the definition of $\alpha_{i}$, it should be orthogonal to every $e_{i}(i>1)$, implying that $\Lambda_{1}$ is proportional to $e_{1}$. Its normalization requires that $\left\langle\alpha_{1}, e_{1}\right\rangle=1$ and since $\alpha_{1}=e_{1}-e_{2}$, we thus have $\Lambda_{1}=e_{1}$. Now, by the same reasoning, $\Lambda_{2}$ should be in the plane spanned by $e_{1}$ and $e_{2}$. Since $\Lambda_{2}$ should be
orthogonal to $\alpha_{1}=e_{1}-e_{2}, \Lambda_{2}$ is proportional to $e_{1}+e_{2}$. In the case $\ell=2$, we thus have $\Lambda_{2}=\frac{1}{2}\left(e_{1}+e_{2}\right)$ while for $\ell>2, \Lambda_{2}=e_{1}+e_{2}$. Assume now (for $i-1<\ell$ ) that From the orthogonality of $\Lambda_{i}$ with $\alpha_{j}(i<j)$ we again deduce that $\Lambda_{i}$ should be a linear combination of $e_{i}(i \leq j)$. Furthermore, the orthogonality of $\Lambda_{i}$ with $\alpha_{j}$ $(j<i)$ implies, based on the relation $e_{i}-e_{j}=\sum_{i \leq k<j} \alpha_{k}, \quad 1 \leq i<j \leq l$, that $\Lambda_{i}$ is orthogonal to $e_{1}-e_{2}, \cdots, e_{1}-e_{i}$ and thus that $\Lambda_{i}$ is proportional to $e_{1}+e_{2}+\cdots+e_{i}$, the coefficient of proportionality being fixed as 1 (resp. 1/2) for $i<\ell$ (resp. $i=\ell$ ). The weight system, and the 8 Weyl chambers, is illustrated in Fig. 3.]
j. Using Weyl formula : $\operatorname{dim}(\Lambda)=\prod_{\alpha>0} \frac{\langle\Lambda+\rho, \alpha\rangle}{\langle\rho, \alpha\rangle}$ compute the dimension of the two fundamental representations of $B_{2}$ and of that of highest weight $2 \Lambda_{2}$. In view of these dimensions, what are these representations of $\mathrm{SO}(5)$ ?
[Solution : In the case of $B_{2}, \Lambda_{1}=e_{1}, \Lambda_{2}=\frac{1}{2}\left(e_{1}+e_{2}\right), \rho=3 e_{1}+e_{2}$. One can check that $\rho=\Lambda_{1}+\Lambda_{2}$. The set of positive roots is $\Delta_{+}=\left\{e_{1}, e_{2}, e_{1} \pm e_{2}\right\}$. Since

$$
\begin{aligned}
\Lambda_{1}+\rho & =\frac{5}{2} e_{1}+\frac{1}{2} e_{2}, \\
\Lambda_{2}+\rho & =2 e_{1}+e_{2}, \\
2 \Lambda_{2}+\rho & =\frac{5}{2} e_{1}+\frac{3}{2} e_{2},
\end{aligned}
$$

we get

$$
\begin{array}{rlrlr}
\left\langle\Lambda_{1}+\rho, e_{1}\right\rangle & =\frac{5}{2}, & \left\langle\Lambda_{1}+\rho, e_{2}\right\rangle & = & \frac{1}{2} \\
\left\langle\Lambda_{1}+\rho, e_{1}-e_{2}\right\rangle & =\frac{5}{2}-\frac{1}{2}=2, & \left\langle\Lambda_{1}+\rho, e_{1}+e_{2}\right\rangle & =\frac{5}{2}+\frac{1}{2}=3, \\
\left\langle\rho, e_{1}\right\rangle & = & \frac{3}{2}, & \left\langle\rho, e_{2}\right\rangle & = \\
\frac{1}{2} \\
\left\langle\rho, e_{1}-e_{2}\right\rangle & =\frac{3}{2}-\frac{1}{2}=1, & \left\langle\rho, e_{1}+e_{2}\right\rangle & =\frac{3}{2}+\frac{1}{2}=2,
\end{array}
$$

and thus

$$
\operatorname{dim}\left(\Lambda_{1}\right)=\prod_{\alpha>0} \frac{\left\langle\Lambda_{1}+\rho, \alpha\right\rangle}{\langle\rho, \alpha\rangle}=\frac{\frac{5}{2} \times \frac{1}{2} \times 2 \times 3}{\frac{3}{2} \times \frac{1}{2} \times 1 \times 2}
$$

i.e.

$$
\operatorname{dim}\left(\Lambda_{1}\right)=5 \quad \text { vector representation. }
$$

See Fig. 4 for the corresponding weight diagram. Similarly, we deduce from

$$
\begin{array}{rlrlrl}
\left\langle\Lambda_{2}+\rho, e_{1}\right\rangle & = & 2, & \left\langle\Lambda_{2}+\rho, e_{2}\right\rangle & =1 \\
\left\langle\Lambda_{2}+\rho, e_{1}-e_{2}\right\rangle & =2-1=1=2, & \left\langle\Lambda_{2}+\rho, e_{1}+e_{2}\right\rangle & =2+1=2
\end{array}
$$



Figure 4 - The weight diagram for the representation $\left(\Lambda_{1}\right)$ of the algebra $B_{2}$. The weights are shown as green dots.
that

$$
\operatorname{dim}\left(\Lambda_{2}\right)=\prod_{\alpha>0} \frac{\left\langle\Lambda_{2}+\rho, \alpha\right\rangle}{\langle\rho, \alpha\rangle}=\frac{2 \times 1 \times 1 \times 3}{\frac{3}{2} \times \frac{1}{2} \times 1 \times 2}
$$

i.e.

$$
\begin{array}{|l|}
\hline \operatorname{dim}\left(\Lambda_{2}\right)=4
\end{array} \text { spinorial representation . }
$$

See Fig. 5 for the corresponding weight diagram. Finally,


Figure 5 - The weight diagram for the representation $\left(\Lambda_{2}\right)$ of the algebra $B_{2}$. The weights are shown as green dots.

$$
\begin{aligned}
& \left\langle 2 \Lambda_{2}+\rho, e_{1}\right\rangle=\frac{5}{2}, \\
& \left\langle 2 \Lambda_{2}+\rho, e_{1}-e_{2}\right\rangle=\frac{5}{2}-\frac{3}{2}=1=2, \quad\left\langle 2 \Lambda_{2}+\rho, e_{1}+e_{2}\right\rangle=\frac{5}{2}+\frac{3}{2}=4,
\end{aligned}
$$

leads to

$$
\operatorname{dim}\left(2 \Lambda_{2}\right)=\prod_{\alpha>0} \frac{\left\langle 2 \Lambda_{2}+\rho, \alpha\right\rangle}{\langle\rho, \alpha\rangle}=\frac{\frac{5}{2} \times \frac{3}{2} \times 1 \times 4}{\frac{3}{2} \times \frac{1}{2} \times 1 \times 2}
$$

i.e.

$$
\operatorname{dim}\left(2 \Lambda_{2}\right)=10 \quad \text { adjoint representation. }
$$

See Fig. 6 for the corresponding weight diagram. Note that $2 \Lambda_{2}=\alpha_{1}+2 \alpha_{2}$ is the


Figure 6 - The weight diagram for the representation $\left(2 \Lambda_{2}\right)$ of the algebra $B_{2}$. The weights are shown as green dots.
heighest root. In the $\alpha_{i}$ basis, it reads $2 \Lambda_{2}=\alpha_{1}+2 \alpha_{2}$ which sum of components is maximal. It is the heighest weight of the adjoint representation. ]
k. Draw on the same figure the roots and the low lying weights of so(5).
[Solution : See Fig. 3.]
2. $\underline{G_{2}}$

In the space $\mathbb{R}^{2}$, we consider three vectors $e_{1}, e_{2}, e_{3}$ of vanishing sum, $\left\langle e_{i}, e_{j}\right\rangle=$ $\delta_{i j}-\frac{1}{3}$, and construct the 12 vectors
$\pm\left(e_{1}-e_{2}\right), \pm\left(e_{1}-e_{3}\right), \pm\left(e_{2}-e_{3}\right), \pm\left(2 e_{1}-e_{2}-e_{3}\right), \pm\left(2 e_{2}-e_{1}-e_{3}\right), \pm\left(2 e_{3}-e_{1}-e_{2}\right)$
They make the root system of $G_{2}$, as we shall check.
a. What can be said on the dimension of the algebra $G_{2}$ ?
[Solution : $\operatorname{dim} G_{2}=$ rank + number of roots $=12+2=14$.]
b. Show that $\alpha_{1}=e_{1}-e_{2}$ and $\alpha_{2}=-2 e_{1}+e_{2}+e_{3}$ are two simple roots, in accord with the Dynkin diagram of $G_{2}$ given in the notes. Compute the Cartan matrix.
[Solution : First,

$$
\left\langle e_{i}, e_{i}\right\rangle=\frac{2}{3} \quad \text { and } \quad\left\langle e_{i}, e_{i}\right\rangle=-\frac{1}{3},
$$

which leads to
$\left\langle\alpha_{1}, \alpha_{1}\right\rangle=2 \frac{2}{3}-2\left(-\frac{1}{3}\right)=2$ and $\left\langle\alpha_{2}, \alpha_{2}\right\rangle=4 \frac{2}{3}+2 \frac{2}{3}-4\left(-\frac{1}{3}\right)-4\left(-\frac{1}{3}\right)+2\left(-\frac{1}{3}\right)=6$,
and

$$
\begin{aligned}
\left\langle\alpha_{1}, \alpha_{2}\right\rangle & =\left\langle\alpha_{2}, \alpha_{1}\right\rangle=-2\left\langle e_{1}, e_{1}\right\rangle+\left\langle e_{1}, e_{2}\right\rangle+\left\langle e_{1}, e_{3}\right\rangle+2\left\langle e_{1}, e_{2}\right\rangle-\left\langle e_{2}, e_{2}\right\rangle-\left\langle e_{2}, e_{3}\right\rangle \\
& =-2 \frac{2}{3}-\frac{1}{3}-\frac{1}{3}-\frac{2}{3}-\frac{2}{3}+\frac{1}{3}=-3 .
\end{aligned}
$$

The Cartan matrix $C_{i j}=2 \frac{\left\langle\alpha_{i}, \alpha_{j}\right\rangle}{\left\langle\alpha_{j}, \alpha_{j}\right\rangle}$ thus reads

$$
C=\left(\begin{array}{cc}
2 & -1 \\
-3 & 2
\end{array}\right)
$$

as expected, in accordance to the Dynkin diagram of Fig. 7.


Figure 7 - The Dynkin diagram for the algebra $G_{2}$.
]
c. What are the positive roots? Compute the vector $\rho$, half-sum of positive roots.
[Solution: $\Delta_{+}=\left\{\alpha_{1}, \alpha_{2}, \alpha_{1}+\alpha_{2}, 2 \alpha_{1}+\alpha_{2}, 3 \alpha_{1}+\alpha_{2}, 3 \alpha_{1}+2 \alpha_{2}\right\}$ and $\rho=5 \alpha_{1}+3 \alpha_{2}$. The root diagram is shown in Fig. 8.]


Figure 8 - The roots of $G_{2}$. In red, the two other simple roots $\alpha_{1}$ and $\alpha_{2}$. In black, the four other positive roots, and in cyan the six negative roots.
d. What is the group of invariance of the root diagram? Show that it is of order 12 and that it is the Weyl group of $G_{2}$. Draw the first Weyl chamber.
[Solution : The group of invariance of the root diagram is the diedral group $D_{6}$, of order 12. The weight diagram is illustrated in Fig. 9]


Figure 9 - The weight diagram for the algebra $G_{2}$. The two simple roots $\alpha_{1}$ and $\alpha_{2}$ are displayed as red lines. The two fundamental weights $\Lambda_{1}$ and $\Lambda_{2}$ are displayed as blue lines. The lattice of roots is shown as red dots, while the lattice of weights is shown as blue dots. These two lattices are identical for $G_{2}$. The 12 Weyl chambers (the first Weyl chambers $\mathcal{C}_{1}$ is explicitely indicated) are shown as grey regions of various intensities. They are obtained by acting on $\mathcal{C}_{1}$ with the various element of the Weyl group.
e. Check that the fundamental weights are

$$
\Lambda_{1}=2 \alpha_{1}+\alpha_{2} \quad \Lambda_{2}=3 \alpha_{1}+2 \alpha_{2} .
$$

[Solution : First, we compute

$$
\left\langle\alpha_{1}, \alpha_{1}\right\rangle=2, \quad\left\langle\alpha_{2}, \alpha_{2}\right\rangle=6 \quad \text { and } \quad\left\langle\alpha_{1}, \alpha_{2}\right\rangle=\left\langle\alpha_{2}, \alpha_{1}\right\rangle=-3 .
$$

Since $\Lambda_{1}$ is orthogonal to $\alpha_{2}$, obviously $\Lambda_{1}$ should be proportional to $2 \alpha_{1}+\alpha_{2}$. The normalization is fixed by the condition $2 \frac{\left\langle\Lambda_{1}, \alpha_{1}\right\rangle}{\left\langle\alpha_{1}, \alpha_{1}\right\rangle}=1$, leading to $\Lambda_{1}=2 \alpha_{1}+$ $\alpha_{2}$. Similarly, $\Lambda_{2}$ is orthogonal to $\alpha_{1}$, and obviously $\Lambda_{2}$ should be proportional to $\alpha_{2}+\frac{3}{2} \alpha 1$. The normalization is fixed by the condition $2 \frac{\left\langle\Lambda_{2}, \alpha_{2}\right\rangle}{\left\langle\alpha_{2}, \alpha_{2}\right\rangle}=1$, leading to $\left.\Lambda_{2}=3 \alpha_{1}+2 \alpha_{2}.\right]$
f. What are the dimensions of the fundamental representations?

## [Solution :

$$
\begin{aligned}
& \left\langle\rho, \alpha_{1}\right\rangle=5\left\langle\alpha_{1}, \alpha_{1}\right\rangle+3\left\langle\alpha_{1}, \alpha_{2}\right\rangle=10-9=1 \\
& \left\langle\rho, \alpha_{2}\right\rangle=5\left\langle\alpha_{1}, \alpha_{2}\right\rangle+3\left\langle\alpha_{2}, \alpha_{2}\right\rangle=-15+18=3
\end{aligned}
$$

and thus

$$
\left\langle\rho, \alpha_{1}+\alpha_{2}\right\rangle=4,\left\langle\rho, 2 \alpha_{1}+\alpha_{2}\right\rangle=5,\left\langle\rho, 3 \alpha_{1}+\alpha_{2}\right\rangle=6,\left\langle\rho, 3 \alpha_{1}+2 \alpha_{2}\right\rangle=9 .
$$

We also easily get

$$
\begin{aligned}
& \left\langle\Lambda_{1}, \alpha_{1}\right\rangle=2\left\langle\alpha_{1}, \alpha_{1}\right\rangle+\left\langle\alpha_{1}, \alpha_{2}\right\rangle=4-3=1 \\
& \left\langle\Lambda_{1}, \alpha_{2}\right\rangle=2\left\langle\alpha_{1}, \alpha_{2}\right\rangle+\left\langle\alpha_{2}, \alpha_{2}\right\rangle=-6+6=0
\end{aligned}
$$

which leads to

$$
\left\langle\Lambda_{1}, \alpha_{1}+\alpha_{2}\right\rangle=1,\left\langle\Lambda_{1}, 2 \alpha_{1}+\alpha_{2}\right\rangle=2,\left\langle\Lambda_{1}, 3 \alpha_{1}+\alpha_{2}\right\rangle=3,\left\langle\Lambda_{1}, 3 \alpha_{1}+2 \alpha_{2}\right\rangle=3 .
$$

Thus
$\operatorname{dim}\left(\Lambda_{1}\right)=\prod_{\alpha>0} \frac{\left\langle\Lambda_{1}+\rho, \alpha\right\rangle}{\langle\rho, \alpha\rangle}=\left(1+\frac{1}{1}\right)(1+0)\left(1+\frac{1}{4}\right)\left(1+\frac{2}{5}\right)\left(1+\frac{3}{6}\right)\left(1+\frac{3}{9}\right)$
and thus

$$
\operatorname{dim}\left(\Lambda_{1}\right)=7,
$$

see Fig. 10 for the corresponding weight diagram. Finally, we get

$$
\begin{aligned}
& \left\langle\Lambda_{2}, \alpha_{1}\right\rangle=3\left\langle\alpha_{1}, \alpha_{1}\right\rangle+2\left\langle\alpha_{1}, \alpha_{2}\right\rangle=6-6=0 \\
& \left\langle\Lambda_{2}, \alpha_{2}\right\rangle=3\left\langle\alpha_{1}, \alpha_{2}\right\rangle+2\left\langle\alpha_{2}, \alpha_{2}\right\rangle=-9+12=3
\end{aligned}
$$

which leads to

$$
\left\langle\Lambda_{2}, \alpha_{1}+\alpha_{2}\right\rangle=3,\left\langle\Lambda_{2}, 2 \alpha_{1}+\alpha_{2}\right\rangle=3,\left\langle\Lambda_{2}, 3 \alpha_{1}+\alpha_{2}\right\rangle=3,\left\langle\Lambda_{2}, 3 \alpha_{1}+2 \alpha_{2}\right\rangle=6 .
$$

We obtain
$\operatorname{dim}\left(\Lambda_{2}\right)=\prod_{\alpha>0} \frac{\left\langle\Lambda_{2}+\rho, \alpha\right\rangle}{\langle\rho, \alpha\rangle}=(1+0)\left(1+\frac{3}{3}\right)\left(1+\frac{3}{4}\right)\left(1+\frac{3}{5}\right)\left(1+\frac{3}{6}\right)\left(1+\frac{6}{9}\right)$


Figure 10 - The weight diagram for the representation $\left(\Lambda_{1}\right)$ of the algebra $G_{2}$. The weights are shown as green dots.


Figure 11 - The weight diagram for the representation $\left(\Lambda_{2}\right)$ of the algebra $G_{2}$. The weights are shown as green dots.
and thus, as expected,

$$
\operatorname{dim}\left(\Lambda_{2}\right)=14 \quad \text { adjoint representation, }
$$

see Fig. 11 for the corresponding weight diagram. Note that again, $\Lambda_{2}=3 \alpha_{1}+2 \alpha_{2}$ is the heighest root.]
g. In the two cases of $B_{2}$ and $G_{2}$, one observes that the highest weight of the adjoint representation is given by the highest root. Explain why this is true in general.
[Solution : The roots are the weights of the adjoint representation. The highest weight of the adjoint representation is thus the highest root.]
3. A little touch of particle physics

Why were the groups $S O(5)$ or $G_{2}$ inappropriate as symmetry groups extending the $\mathrm{SU}(2)$ isospin group, knowing that several "octets" of particles had been observed?
[Solution : There is no representation of dimension 8, although $7+1$ would be not so bad ... ?]
D. Dimensions of $S U(3)$ representations

We admit that the construction of § 3.4.2, of completely symmetric traceless rank $(p, m)$ tensors in $\mathbb{C}^{3}$, does give the irreducible representations of $\mathrm{SU}(3)$ of highest weight $(p, m)$. Then we want to determine the dimension $d(p, m)$ of the space of these tensors.

1. Show, by studying of the product of two tensors of rank $(p, 0)$ and $(0, m)$ and separating the trace terms (those containing a $\delta_{j}^{i}$ between one lower and one upper index) that $(p, 0) \otimes(0, m)=((p-1,0) \otimes(0, m-1)) \oplus(p, m)$ and thus that

$$
d(p, m)=d(p, 0) d(0, m)-d(p-1,0) d(0, m-1) .
$$

[Solution : The tensors of the space $(p, 0) \otimes(0, m)$ are tensors of rank $(p, m)$ completely symmetric in their $p$ upper indices, completely symmetric in their $m$ lower indices, but they have a priori arbitrary traces between upper and lower indices. We would like to show that any tensor $t_{j_{1} \cdots j_{m}}^{i_{1} \cdots i_{p}}$ of this space can be expressed as the sum
of a tensor with the same symmetries and with vanishing traces, and of a tensor which possesses traces, i.e. of the form

$$
v_{j_{1} \cdots j_{m}}^{i_{1} \cdots i_{p}}:=\sum_{n=1}^{m} \sum_{q=1}^{p} \delta_{j_{n}}^{i_{q}} u_{j_{1} \cdots \hat{j}_{n} \cdots j_{m}}^{i_{1} \cdots \hat{q}_{q} \cdots i_{p}}
$$

where the hat on top of an index means that the given index has been omitted, and $u$ is a tensor to be determined, which is completely symmetric in its $p-1$ upper indices and in its $m-1$ lower indices, thus belonging to the space $(p-1,0) \otimes(0, m-1)$. One would thus like to write

$$
t_{j_{1} \cdots j_{m}}^{i_{1} \cdots i_{p}}=[t-v]_{j_{1} \cdots j_{m}}^{i_{1} \cdots i_{p}}+v_{j_{1} \cdots j_{m}}^{i_{1} \cdots i_{p}},
$$

$[t-v]$ being traceless, and we are going to fix $u$ by requiring that $\delta_{i_{1}}^{j_{1}}[t-v]_{j_{1} \cdots j_{m}}^{i_{1} \cdots i_{p}}=0$. (due to the symmetries of $t$ and $v$ this implies that every trace betwen an upper index and a lower index vanish). We first consider the case $p=m=2$, and write

$$
v_{j_{1} j_{2}}^{i_{1} i_{2}}=\delta_{j_{1}}^{i_{1}} u_{j_{2}}^{i_{2}}+\delta_{j_{1}}^{i_{2}} u_{j_{2}}^{i_{1}}+\delta_{j_{2}}^{i_{1}} u_{j_{1}}^{i_{2}}+\delta_{j_{2}}^{i_{2}} u_{j_{1}}^{i_{1}} .
$$

Since each index can take 3 values, the constraint

$$
\delta_{i_{1}}^{j_{1}}[t-v]_{j_{1} j_{2}}^{i_{1} i_{2}}=0
$$

leads to

$$
t_{i j_{2}}^{i i_{2}}-\left[3 u_{j_{2}}^{i_{2}}+u_{j_{2}}^{i_{2}}+u_{j_{2}}^{i_{2}}+\delta_{i_{2}}^{j_{2}} u_{i}^{i}\right]=0 .
$$

Taking now a new trace by mean of $\delta_{i_{2}}^{j_{2}}$, we get $8 u_{i}^{i}=t_{i j}^{i j}$. This completely fixes the expression of $u_{j_{2}}^{i_{2}}$ according to

$$
5 u_{j_{2}}^{i_{2}}=t_{i j_{2}}^{i i_{2}}-\delta_{i 2}^{j_{2}} u_{i}^{i}=t_{i j_{2}}^{i i_{2}}-\frac{1}{8} \delta_{i 2}^{j_{2}} t_{i j}^{i j},
$$

i.e.

$$
u_{j_{2}}^{i_{2}}=\frac{1}{5} t_{i j_{2}}^{i i_{2}}-\frac{1}{40} \delta_{i_{2}}^{j_{2}} t_{i j}^{i j} .
$$

One can similarly study the case $p=3$ and $m=2$, and write

$$
v_{j_{1} j_{2}}^{i_{1} i_{2} i_{3}}=\delta_{j_{1}}^{i_{1}} u_{j_{2}}^{i_{2} i_{3}}+\delta_{j_{1}}^{i_{2}} u_{j_{2}}^{i_{1} i_{3}}+\delta_{j_{1}}^{i_{3}} u_{j_{2}}^{i_{1} i_{2}}+\delta_{j_{2}}^{i_{1}} u_{j_{1}}^{i_{2} i_{3}}+\delta_{j_{2}}^{i_{2}} u_{j_{1}}^{i_{1} i_{3}}+\delta_{j_{2}}^{i_{3}} u_{j_{1}}^{i_{1} i_{2}} .
$$

The contraction of $\delta_{i_{1}}^{j_{1}}$ with $[t-v]_{j_{1} j_{2}}^{i_{1} i_{2} i_{3}}$ gives

$$
[3+1+1+1] u_{j_{2}}^{i_{2} i_{3}}+\delta_{j_{2}}^{i_{2}} i_{i}^{i i_{3}}+\delta_{j_{2}}^{i_{3}} u_{i}^{i i_{2}}-t_{i j_{2}}^{i i_{2} i_{3}}=0 .
$$

Next, contracting with $\delta_{j_{2}}^{i_{2}}$ gives

$$
[6+3+1] u_{i}^{i i_{3}}-t_{i j}^{i j i_{3}}=0,
$$

and thus allows us to fix completely $u$ as

$$
u_{j_{2}}^{i_{2} i_{3}}=-\frac{1}{60}\left[t_{i j}^{i j i_{3}} \delta_{j_{2}}^{i_{2}}+t_{i j}^{i j i_{2}} \delta_{j_{2}}^{i_{3}}\right]+\frac{1}{6} t_{i j_{2}}^{i i_{2} i_{3}} .
$$

The general case can be treated following the same line of thought, although it is rather painful to write explicitely : after a finite number of operation (equal to $\inf (p, m)$ ), one can determine completely the tensor $u$, which achieves the proof that $(p, 0) \otimes(0, m)=((p-1,0) \otimes(0, m-1)) \oplus(p, m)$ and thus that

$$
d(p, m)=d(p, 0) d(0, m)-d(p-1,0) d(0, m-1) .
$$

]
2. Show by a computation analogous to that of $\mathrm{SU}(2)$ that

$$
d(p, 0)=d(0, p)=\frac{1}{2}(p+1)(p+2) .
$$

[Solution : By symmetry, one can always organize the indices in such a way that 1's occurs first, then 2's and finally 3's, i.e. looking as $1 \cdots 12 \cdots 23 \cdots 3$. On should niw introduce 2 separations between the blocks of 1,2 and 3 's. Consider first the separation between 2 and 3's. There are $p+1$ possibilities (from a configuration with only 3 's, looking as $3 \cdots 3$, till the one with no 3's, looking as $1 \cdots 12 \cdots 2$ ), labeling for example the position $r$ of the first 3 (which varies from 1 to $p+1$ in these two extreme cases). Then, the total number of 1's and 2's is $r-1$, and we thus have $r$ possibilities to choose the separation between 1's and 2's. This means that the number we are looking for is just

$$
d(p, 0)=d(0, p)=\sum_{r=1}^{p+1} r=\frac{1}{2}(p+1)(p+2) .
$$

]
3. Derive from it the expression of $d(p, m)$ and compare with (3.64).
[Solution: From $d(p, m)=d(p, 0) d(m, 0)-d(p-1,0) d(m-1,0)=\frac{1}{4}(p+1)(p+$ $2)(m+1)(m+2)-\frac{1}{4} p(p+1) m(m+1)$ we deduce

$$
d(p, m)=\frac{1}{2}(m+1)(p+1)(m+p+2),
$$

